

# Lightlike form factors and OPE

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Based on the work with Yuanhong Guo, Lei Wang  
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# Outline:

- I. Introduction
- II. Bootstrap  $\sim$  2-loop  $F_4^{(2)}$
- III. FFOPE  $\sim$  Non-perturbative
- IV. Summary & Outlook.

# I. Introduction and some history.

What are form factors?

EM vertex:

$$\Gamma^\mu = \int \bar{\psi} \gamma^\mu \psi = \text{tree} + \text{loop} + \dots$$

$$= \gamma^\mu \underbrace{F_1(q^2)}_{\text{Form factor}} + i \frac{\sigma^{\mu\nu} q_\nu}{2m} \underbrace{F_2(q^2)}_{\text{Form factor}}$$

$$g-2 = 2F_2(q^2=0) = \frac{\alpha}{\pi} + \mathcal{O}(\alpha^2)$$

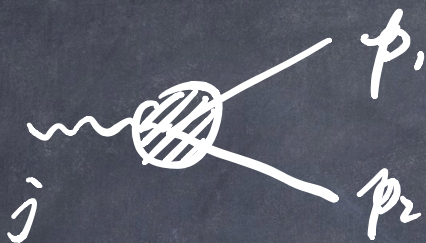


$\sigma \sim \text{depend on } \{F_1, F_2\}$

Nuclear form factor:  $F_1, F_2$

↳ characterize shape of particle  
deviation from point particle.

# Sudakov Form factor.



- Sudakov (1954) sum over leading log.

(off-shell)

$$\Gamma(p_1, p_2, q) \sim \exp \left[ -e^2 \log \left| \frac{q^2}{p_1^2} \right| \log \left| \frac{q^2}{p_2^2} \right| \right]$$

$$q^2 \gg p_1^2, p_2^2 \gg m^2$$

- Jackiw (1966)  $\rightarrow$  PhD thesis high-energy form factor

(on-shell)

$$\Gamma(q) \sim \exp \left[ -e^2 \log^2 \frac{q^2}{\mu^2} \right]$$

Further development  $\sim 1980$ .

- A. Mueller (79')  $\leadsto$  sub-leading log. resum.  
J. Collins (80') in QED.
- A. Sen (81')  $\rightarrow$  Non-abelian.

Factorization



$\sim (\text{soft}) \times (\text{hard})$

$\downarrow$   
Sudakov FF.

"4-Loop."

# General high-point FF.

F.T.

$$\int d^D x e^{-i q x} \langle p_1 p_2 \dots p_n | U(x) | 0 \rangle$$

$$F_{n,U}(p_1 \dots p_n)$$

Asymp. on-shell states.

local oper.



$$q = \sum_{i=1}^n p_i$$

$$q^2 \neq 0$$

$$\langle p_1 \dots p_n | 0 \rangle$$

Amplitudes.

$$\langle U_1 O_2 \dots U_n \rangle$$

correlation funct.

2010 draw attentions.

- Maldacena, Zhibodov. 2010.  
using AdS/CFT study strong coupling FF.
- Brandhuber et.al. 2010. } weak coupling.  
Bork et.al. 2010.

MHV FF:

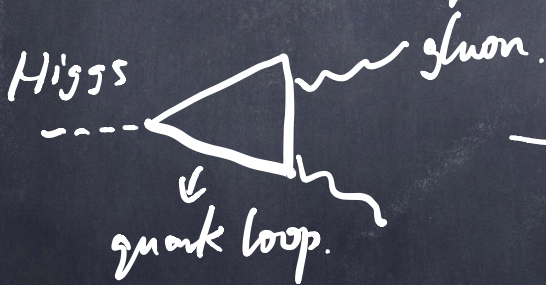
$$F_{\text{tr}(\phi^2)}^{(0), \text{MHV}}(1^g \dots i^\phi \dots j^\phi \dots n^g) = \frac{\langle ij \rangle^2 \delta^4(\sum p - q)}{\langle 12 \rangle \langle 23 \rangle \dots \langle n1 \rangle}$$

$\text{tr}(F^2)(1^{g^+} \dots i^{g^-} \dots j^{g^-} \dots n^{g^+})$   $\rightarrow$  apply to  $N=4$  QCD



# Applications of FF:

- Sudakov FF (IR divergence)  
↳ in general gauge theory.
- Anomalous dimensions (UV renormalization)
- EFT Amplitudes.



integrate over quarks.



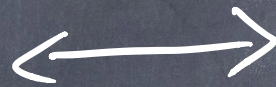
$$\underline{H \text{tr}(F_{\mu\nu} F^{\mu\nu})} + \text{higher dim-vert.}$$

Higgs + gluons Amp  $\leftrightarrow$  Formfactor with  $\text{tr}(F^2)$   
 $q^2 = m_H^2$



$F_{\text{tr}(F^2)}^{(2)}(123)$  in  $N=4$

Brandhuber et al. 2012.



in QCD.

Gehrmann et al.  
2011.

Maximal Transcendental

Principle of Max. Transcendentality.

Lipatov et al. 2001.

$$\gamma^{N=4} = \gamma^{QCD} \Big|_{L.T.}$$

This talk we focus on

"Lightlike" form factor. in  $N=4$  SYM.

$$F_{4.}^{\mathcal{L}} = \langle 1^g 2^g 3^g 4^g | \mathcal{L}(q) | 0 \rangle$$

$\rightarrow$   $\text{tr}(F^2)$  + super comp.

• Bootstrap to 2-loops.

$$q^2 = 0$$

• FF OPE.

$$q = \sum_i p_i$$

# Bootstrap Method.

$$\text{Loop Amp/FF} = \sum_i \text{Coeff}_i \times \text{Master-Integral}_i$$

Physical information.

Theory-independent  
Known in many  
cases.  
(once for all).

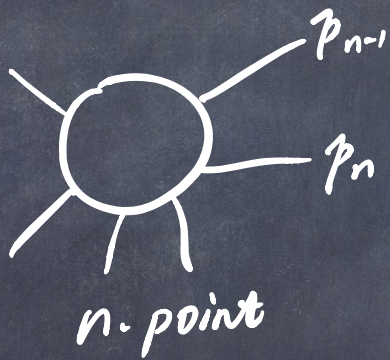
- Traditional method. "Top-down".
- Bootstrap "bottom-up".

Use physical constraints directly to fix final result.

# Constraints :

- IR divergence:  $\frac{1}{\epsilon^\#}$  divergent terms are <sup>well-</sup>known.

- Collinear limit :



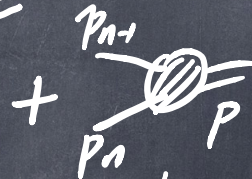
$p_{n-1} \parallel p_n$



(n-1) point

↓  
Lower point

well-known.



↓  
splitting factor.

Universal.

BDS straction : Bern, Dixon, Smirnov 2005

$$\underline{F^{(1)}} , \underline{F^{(2)}}$$

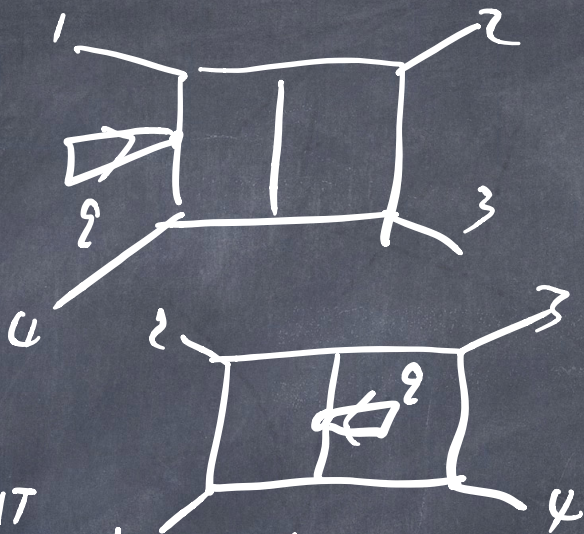
define Remainder function:

$$R^{(2)} = \underline{F_{(\epsilon)}^{(2)}} + \frac{1}{2} \left( F_{(\epsilon)}^{(1)} \right)^2 - \underline{f_{(\epsilon)}^{(2)}} * F_{(2\epsilon)}^{(1)}$$

$$\underline{f^{(2)}(\epsilon)} = -2 \zeta_2 - 2 \zeta_3 \epsilon - 2 \zeta_4 \epsilon^2 \rightarrow \text{contains splitting info.}$$

$R^{(2)}$  : ① is finite.

② has trivial collinear limit  $R_n^{(2)} \xrightarrow{P_n/P_{n-1}} R_{n-1}^{(2)}$



① Ansatz

$$F_4^{(2)} = \sum_{i=1}^{590} C_i \underbrace{I_i^{(2), UT}}_{\text{un-known.}}$$

$2^2 = 0$  all master integral are known.

② One-loop is known

③ Apply constraints.



$$\frac{\#}{\epsilon^4} + \frac{\#}{\epsilon^3} + \dots$$

\* Symmetry  $D_4$

$$F_4^{(2)} = F_4^{(2)} \Big|_{p_i \rightarrow p_{i+1}} = F_4^{(2)} \Big|_{p_i \rightarrow 5-p_i}$$

\* IR : 
$$\underline{F_4^{(2)}} \stackrel{\text{IR}}{\sim} \frac{1}{2} (F_4^{(1)})^2 + f^{(2)} F_{(2\epsilon)}^{(1)}$$

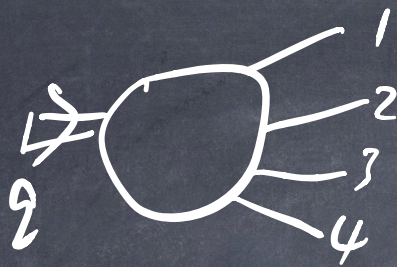
\* Collinear.

$$R_4^{(2)} \xrightarrow{p_i \neq p_{i+1}} \underline{R_3^{(2)} = \text{const.}}$$

Known.

$\Rightarrow$  We can fix  $R_4^{(2)}$  uniquely!  $\checkmark$





$$p_1, p_2, p_3, p_4, \quad q = \sum_{i=1}^4 p_i$$

$$S_{ij} : S_{12}, S_{13}, S_{14}, S_{23}, S_{24}, S_{34}$$

$$\bullet \quad q^2 = 0 \Rightarrow \sum_{i < j} S_{ij} = 0 \Rightarrow 5\text{-independent}$$

$$\bullet \quad \text{scale invariance. } R_4^{(2)} \sim \frac{S_{ij}}{S_{12}} \Rightarrow 4\text{-independent}$$

Result shows that there are only

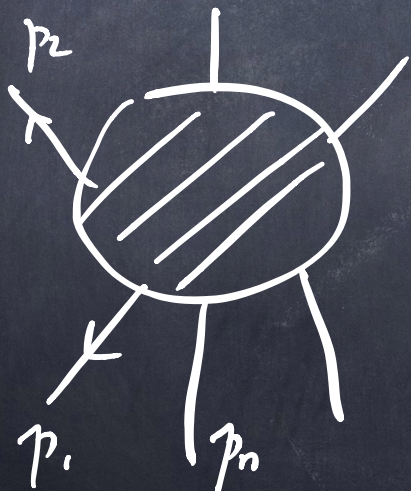
3 independent ratio variables!

→ imply there is a hidden symmetry!

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Duality between Amp/FF and Wilson Loop.

Alday-Maldacena 07. (strong coupling)



Amp.



$$x_{i+1} - x_i = p$$

Wilson-loop

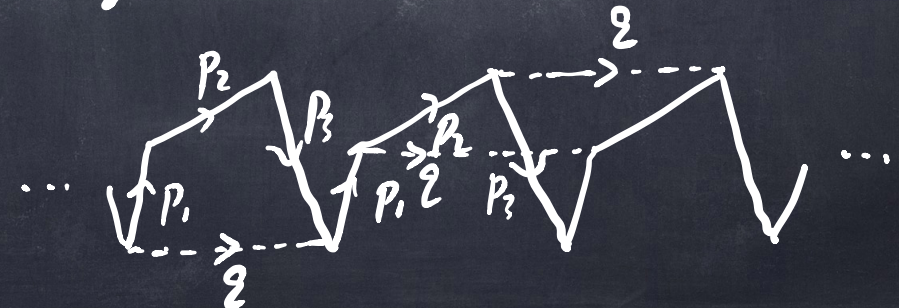
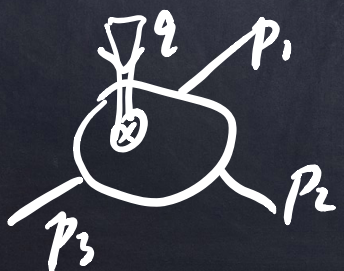
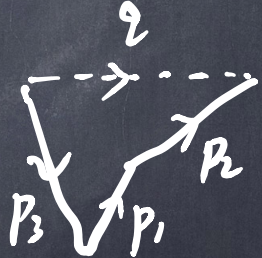
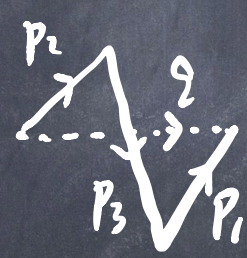
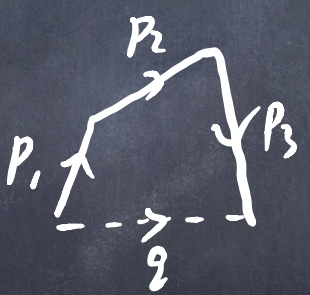
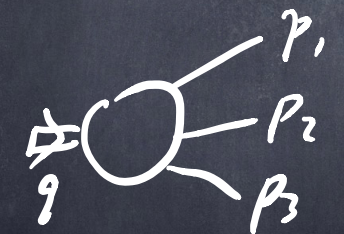
Dual  
conf sym.  
Hidden sym!

$$KW(C) = \int \text{tr} P e^{\int A_\mu(x) dx}$$

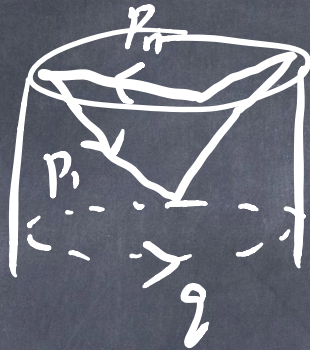
$\downarrow$  conf sym  
 $\downarrow$  Dual conf sym

Form factor:

3-pt example:



form factor  $\leftrightarrow$  periodic Wilson loop



one  
 $q \sim$  period



$$q^2 \neq 0$$

non-lightlike

C.T.

straightline



curved line

$$q^2 = 0$$

Lightlike

C.T.



Lightlike.

↙ A dual conformal sym along  $g$ -direction:

$$\delta_g x_i^m = \frac{1}{2} x_i^2 g^m - (x_i \cdot l) x_i^m$$

$$\delta_g \langle W.L. \rangle = 0 = \underline{\underline{\delta_g R^{(L)}}}$$

↳ one sym  $\rightarrow$  eliminate one variable as mentioned above.

$$A_n, \text{pentagon } W_n : \underline{3n-15}$$

$$F_{n, g^2=0} : 3n-9$$

$$F_{n, g^2 \neq 0} \text{ zigzag } : 3n-7$$

$$n=4 : 3 \times 4 - 7 = 5$$

$$n=4 : 3 \times 4 - 9$$

$$= \underline{\underline{3}}$$

$$\{\mathcal{P}_1, \dots, \mathcal{P}_n\} \quad \sum_i \mathcal{P}_i = 0$$

$$\mathcal{P}_i^2 = 0$$

$$3n - 15$$

$\underbrace{\hspace{2cm}}$   
 conf group  
 $10 + 4 + 1$   
 $K^u$

$$\{\mathcal{P}_1, \dots, \mathcal{P}_n, \mathcal{Q}\} \quad \mathcal{P}_i^2 = 0$$

$$\sum_i \mathcal{P}_i = \mathcal{Q}$$

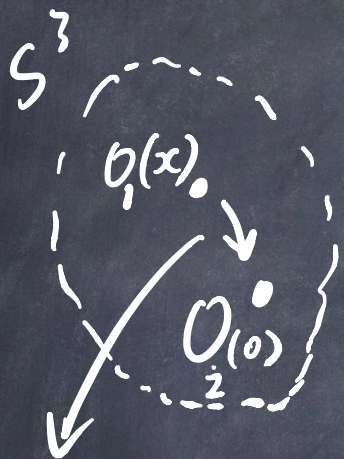
$$3n - (15 - 4 - 4)$$

$\bar{p}^u \quad \bar{K}^u$

$$= 3n - 7$$

OPE

# Analogy with usual OPE in CFT.



Dilatation transformation.

Hamiltonian.

$$O_1(x) O_2(0) \sim \sum C_{12i} \underbrace{f(x, \partial_x)}_{\text{structure const.}} \underbrace{O_i}_{\text{determine by C.S. up to } \Delta, \ell}$$

structure const.

determine by C.S.

up to  $\Delta, \ell$ .

$O_i$  characterized by  $\Delta, \ell$ .  
 $\downarrow$   
 $D O_i = \Delta_i O_i$



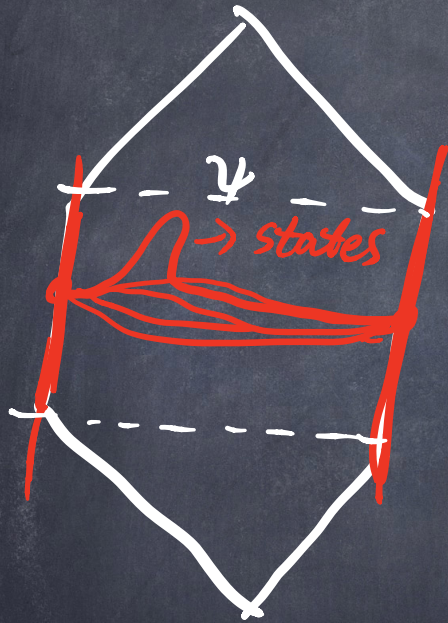
5 pt correlator  $\langle U_1 U_2 \dots U_5 \rangle$

$$\sim \sum_{\psi_1, \psi_2} \frac{C_{12\psi_1}}{C_{45\psi_2}} \cdot \text{conf. PW.}(\Delta_i)$$

$\{ C_{12\psi_1}, \Delta_i \}$  CFT data.



# Wilson Loop OPE.

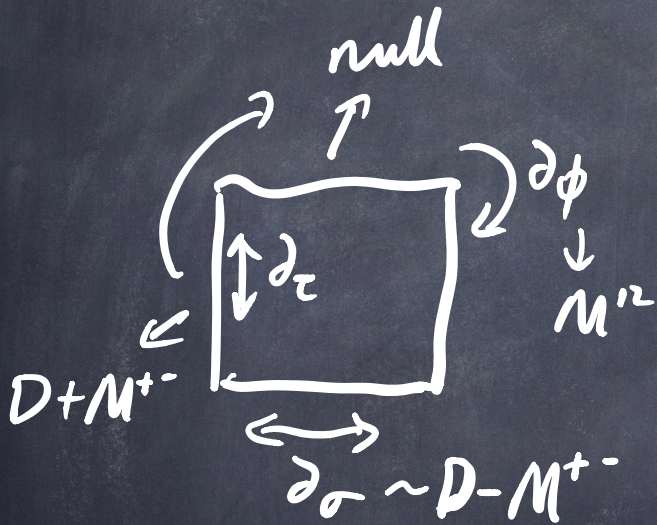


$$W(\square) = \sum_{\psi} e^{-E_{\psi} \tau} \underline{\underline{C_{\psi}}}$$

Alday, Gaiotto, Maldacena,  
Sever, Vieira 2010.

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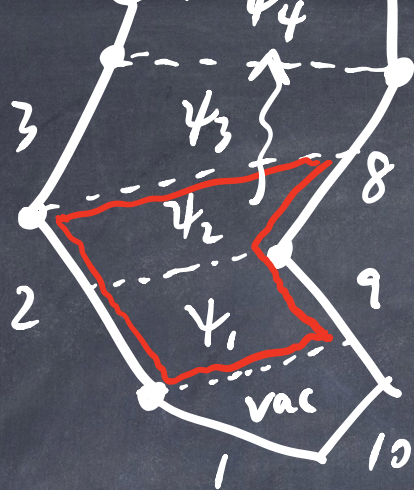
The expansion states are labeled by symmetries of a square null.



$\partial_\tau, \partial_\sigma, \partial_\phi$  对应  $U(1)$   
 $\downarrow \quad \downarrow \quad \downarrow$   
 $E \quad p \quad m$



$\psi_i \sim \{ \sigma_i, \tau_i, \phi_i \}$   
 $\downarrow \# = 5$



$$3n - 15 \Big|_{n=10} = 15 = \underline{\underline{3 \times 5}}$$

$$W_{10} = \sum_{\{\psi_i\}} \prod_{i=1}^5 e^{-\tau_i E_i + i p_i \sigma_i + i m_i \phi_i}$$

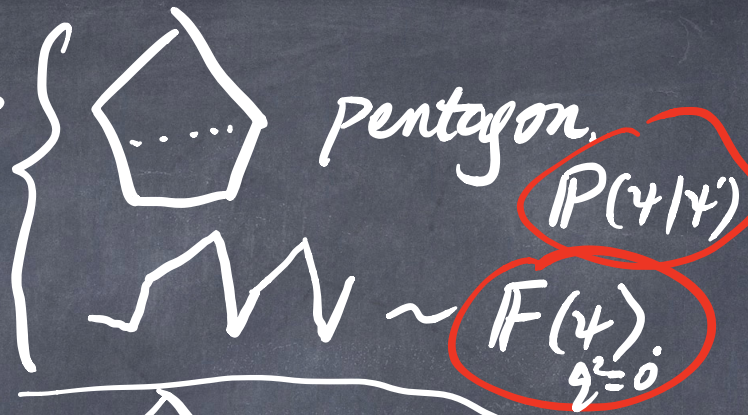
$$\frac{P(\text{vac}|\psi_1) P(\psi_1|\psi_2)}{\dots P(\psi_5|\text{vac})}$$

Analogy :



null box.

$\langle \dots \rangle$  lijk  $\longleftrightarrow$



$\langle :: \rangle$   $G_4$   $\longleftrightarrow$



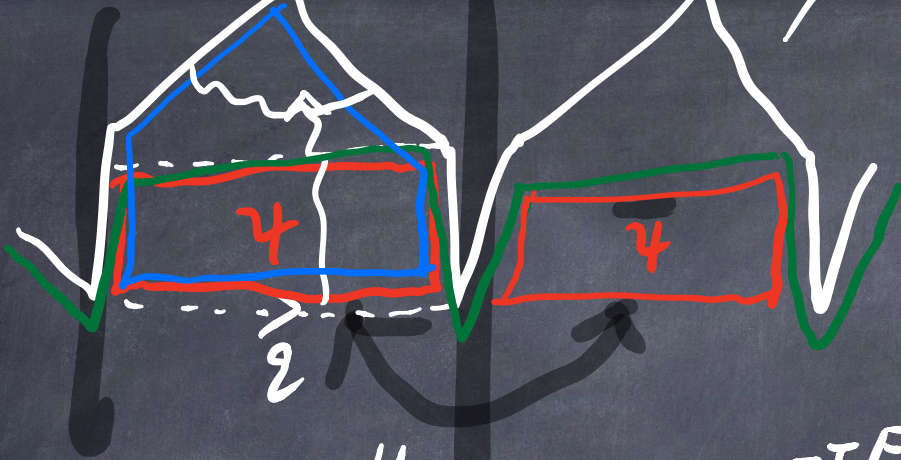
Form factor.

$\underline{3n-9}$   $n=3$  const.



based on  
Sever. Tumanov.  
Wilhelm 2020.

$g^2 \neq 0$



$$\psi \sim (\tau, \phi, \sigma)$$

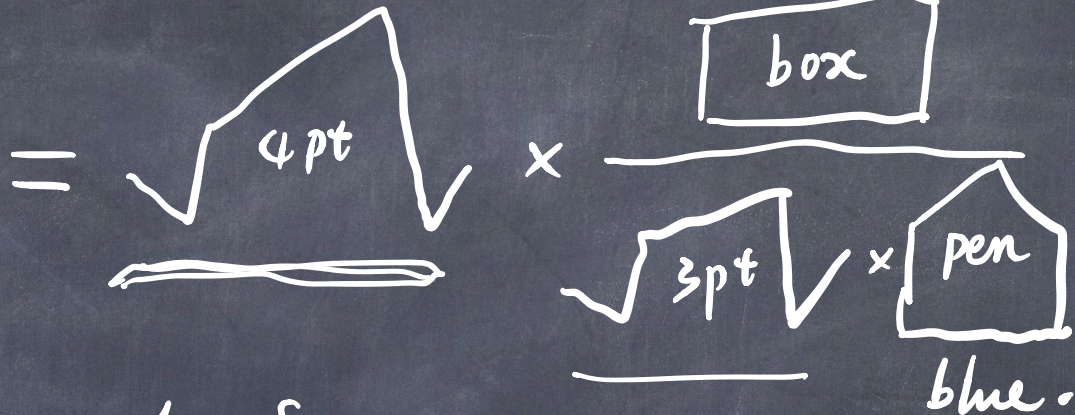
$$\triangle F_4^{\text{LL.}} = \sum e^{-\tau E + i p \sigma + i m \phi}$$

$$\underline{P(\text{vac}|\psi)} \underline{F(\psi)}$$

New building block  
for Lightlike FF.

Regularization:

$$\underline{\underline{W_4^u = F_4^{\text{Reg}}}}$$



Function of  $\{\sigma_i, \tau_i, \phi_i\}$

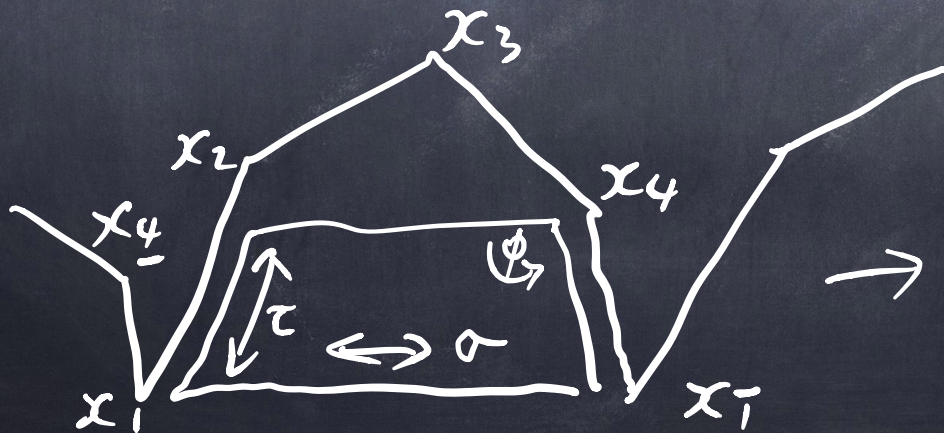
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$$W_4^{u(1)} = -\frac{1}{2} \log\left(\frac{(1-u_1)(1-u_2)}{1-u_3}\right) \log\left(\frac{(1-u_1)(1-u_2)}{1-u_3} \frac{u_3}{u_1}\right)$$

$$u_1 = \frac{S_{12}}{S_{34}} = \frac{\chi_{43}^2}{\chi_{31}^2} = \underline{e^{-2\sigma}}$$

$$u_2 = \frac{S_{23}}{S_{41}} = \frac{\chi_{24}^2}{\chi_{42}^2} = \underline{e^{-2\tau}}$$

$$u_3 = \frac{S_{123} S_{341}}{S_{234} S_{412}} = \frac{\chi_{14}^2 \chi_{32}^2}{\chi_{21}^2 \chi_{43}^2} = \underline{\frac{\cosh(\sigma - \tau) + \cos\phi}{\cosh(\sigma + \tau) + \cos\phi}}$$



$$\delta q u_i = 0$$

From perturbative.: Large  $\tau$  expansion

$$W_4^{(1)} = \underline{-2 e^{-2\tau} \cos^2 \phi} + O(e^{-3\tau})$$

$$W_4^{(2)} = 2 e^{-\tau} \cos \phi h^{(2)}(\sigma)$$

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↓ Map to OPE.

$$W_4^u = 1 + \cos \phi \int \frac{du}{2\pi} e^{-E(u)\tau + iP(u)\sigma} \times \frac{\mathbb{P}(\sigma|u) \overline{\mathbb{F}}_F(u)}{\mu_F u}$$

(first excitation).

$\tau$ -dependence.  
↑



$E(u)$ ,  $P(u)$ ,  $P(0|u)$ ,  $\underline{\underline{\mu_F(u)}}$ .

△ Known in previous work!

→ From <sup>new</sup> perturbative  
we can determine.

$$\mathbb{F}_F(u) = g^0 \times 0 + g^2 \times \underline{\underline{(*)}} + V(g^4)$$

$$\mathbb{E}(u) = \underline{\underline{1}} + \gamma^{(1)} \underbrace{(g^2)}_{-2\mu_F(u)/g^2} + \dots$$

$$\underline{W_4^{(l)}} = \underline{e^{-\tau}} \left( \tau^0 \bar{\Sigma}_0 + \tau^1 \bar{\Sigma}_1 + \dots + \tau^{l-2} \bar{\Sigma}_{l-2} \right)$$

$$\bar{\Sigma}_{l-1} = 0 \quad (\text{from 1-loop})$$

$$\bar{\Sigma}_{l-2} = \int \frac{du}{2\pi} e^{2iu\sigma} \left( \frac{(\hat{\gamma}^{(1)})^{l-2}}{(l-2)!} \right) \times \left[ \mu_{\tilde{r}}(u) \right]_{g^2}^2$$

$\downarrow$   
 $g^2$   
 2 loop data.

$$\underline{e^{-\gamma^{(1)} g^2}} \quad (g^2)^{l-2}$$

This is for any loop order!

# Summary & Outlook.

•  $\downarrow$  Amp/FF  
momenta.



WL ( $\square/\Omega$ )

$\downarrow$   
geometrical  
2D.

• Bootstrap + OPE



